

V. ROSENHAUS, K. KIIRANEN

ON SYMMETRY GROUPS OF YANG-MILLS AND SELF-DUAL YANG-MILLS EQUATIONS

(Presented by P. Kard)

The problem of finding groups of point transformations admissible by Yang-Mills equations and some forms of self-dual Yang-Mills equations in four-dimensional Euclidean space, is considered. The explicit form of basic generators for such groups is found. Some conserved quantities are obtained.

Introduction

It is well known that symmetry plays an important role in physics. There is a connection between symmetry groups and conserved quantities of a differential equation. The theory of group properties of differential equations worked out by S. Lie [1,2], F. Engel [2], A. Tresse [3], E. Vesiot [4], A. Cohen [5], L. V. Ovsyannikov [6] and others, has several applications in mechanics and other branches of physics [7-11].

Here we call the reader's attention to some field equations interesting from the viewpoint of particle physics, such as the Yang-Mills and the so-called self-dual Yang-Mills equations.

Our aim is to find their symmetry groups as well as their trivial conservation laws. This forms a continuation of the investigation of field equations which we started with two-dimensional equations [12] and then carried on with four dimensional ones. Detailed calculations are mostly omitted.

Yang-Mills equations

In 1954 C. N. Yang and R. L. Mills presented [13] the local gauge invariance principle involving an isotopic spin. Compared to electrodynamics, a non-abelian symmetry was introduced.

The Yang-Mills equations are

$$\partial_\mu F_{\mu\nu} - [A_\mu, F_{\mu\nu}] = 0 \quad (1)$$

with the Lagrangian *

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a.$$

Here

$$F_{\mu\nu} = F_{\mu\nu}^a T^a = A_{\mu,\nu} - A_{\nu,\mu} + [A_\mu, A_\nu],$$

$$A_\mu = A_\mu^a T^a, \quad A_{\mu,\nu} \equiv \frac{\partial A_\mu}{\partial x_\nu}.$$

* Here the sum convention is observed. The Greek indices take values from 1 to 4, the Latin ones — from 1 to 3.

The matrices T^a satisfy the conditions

$$[T^a, T^b] = t^{abc} T^c, \\ \text{tr}[T^a, T^b] = -2\delta^{ab},$$

where the structure coefficients t^{abc} are completely antisymmetric.

We study the Yang-Mills equations in the Lorentz gauge:

$$A_{\mu, \mu} = 0, \quad (2)$$

in which equations (1) take the form:

$$\omega_v \equiv T^c [A_{v, \mu \mu}^c + t^{abc} A_{\mu}^a (A_{\mu, v}^b - 2A_{v, \mu}^b + A_{\mu}^k A_v^g t^{kgb})] = 0. \quad (3)$$

We have found the algebra of the point transformation group admissible by the system of differential equations $\omega_v = 0$, following the technique worked out by S. Lie [1-2] and developed further by L. V. Ovsyannikov [6], N. H. Ibragimov [7, 8], S. A. Vladimirov [9] and others. This algebra is presented by generators

$$X = \xi^\mu(x, A) \frac{\partial}{\partial x^\mu} + \eta^\alpha(x, A) \frac{\partial}{\partial A_\alpha^c}.$$

Coefficients ξ^μ and η^α are determined by the equation

$$\tilde{X}\omega_v|_{\omega_v=0} = 0,$$

where

$$\tilde{X} = X + \zeta_\mu^\alpha \frac{\partial}{\partial p_\mu^\alpha} + \sigma_{\mu\nu}^\alpha \frac{\partial}{\partial r_{\mu\nu}^\alpha}$$

is the prolonged generator. The new coefficients ζ_μ^α and $\sigma_{\mu\nu}^\alpha$ depend on the previous ones in the following way:

$$\zeta_\mu^\alpha = D_\mu \eta^\alpha - p_v^\alpha D_\mu \xi^v, \\ \sigma_{\mu\nu}^\alpha = \bar{D}_\mu \zeta_{\nu}^\alpha - r_{\mu\sigma}^\alpha D_\nu \xi^\sigma,$$

where

$$D_\mu \equiv \frac{\partial}{\partial x^\mu} + p_\mu^a \frac{\partial}{\partial A_\alpha^a}, \quad \bar{D}_\mu \equiv D_\mu + r_{\mu\nu}^a \frac{\partial}{\partial p_\nu^a}, \\ p_\mu^a \equiv \frac{\partial A_\alpha^a}{\partial x^\mu}, \quad r_{\mu\nu}^a \equiv \frac{\partial^2 A_\alpha^a}{\partial x^\mu \partial x^\nu}.$$

In the case of Yang-Mills equations, the system of determining equations is the following:

$$\tilde{X}\omega_v|_{(2)(3)} = 0, \\ \tilde{X}p_\mu^\alpha|_{(2)(3)} = 0$$

or, in detail,

$$\begin{aligned} \sigma_{\mu\mu}^c + t^{abc} [\eta^\mu (A_{,\nu}^a - 2A_{,\mu}^b + A^\mu A^\nu t^{hgb}) + \dot{A}^\mu (\zeta_\nu^a - 2\zeta_\mu^b) + \\ + t^{hgb} \dot{A}^\mu (\eta^\mu \dot{A}_\nu^h + \eta^\nu \dot{A}_\mu^h)] |_{(2)(3)} = 0, \\ \zeta_\mu^c |_{(2)(3)} = 0. \end{aligned} \quad (4)$$

First, inserting in (4) identities (3) and $r_{\mu\alpha}^\mu = 0$ (the latter is a consequence of the Lorentz gauge condition (2)) and secondly, equating the coefficients at all remaining independent variables to zero, we get the solutions of the system (4)

$$\xi_\mu = a_{\mu\nu} x_\nu + d \cdot x_\mu + c_\mu, \quad \eta_\mu^c = a_{\mu\nu} \dot{A}_\nu^c - d \cdot \dot{A}_\mu^c + \varepsilon_{ca} \dot{A}_\mu^a,$$

where $a_{\mu\nu} = -a_{\nu\mu}$, $\varepsilon_{ab} = -\varepsilon_{ba}$, c_μ , d are constants. One can see that the generators of the symmetry group of Yang-Mills equations (3) in the Lorentz gauge (2) are:

$$\begin{aligned} X_\mu = \frac{\partial}{\partial x_\mu}, \quad X_{\mu\nu} = Z_{\mu\nu} + Y_{\mu\nu}, \\ X = x_\mu \frac{\partial}{\partial x_\mu} - \dot{A}_\mu^c \frac{\partial}{\partial A_\mu^c}, \quad X_{ab} = \dot{A}_\mu^a \frac{\partial}{\partial A_\mu^b} - \dot{A}_\mu^b \frac{\partial}{\partial A_\mu^a}, \end{aligned} \quad (5)$$

where

$$Z_{\mu\nu} = x_\mu \frac{\partial}{\partial x_\nu} - x_\nu \frac{\partial}{\partial x_\mu}, \quad Y_{\mu\nu} = \dot{A}_\mu^c \frac{\partial}{\partial A_\nu^c} - \dot{A}_\nu^c \frac{\partial}{\partial A_\mu^c}.$$

The rotation generators X_{ab} form a gauge group $SO(3) \simeq SU(2)$. It is to be noted that the Yang-Mills equations without the gauge conditions admit a special conformal transformation

$$X = \left(c_\rho x_\rho x_\mu - \frac{1}{2} c_\mu x_\rho^2 \right) \frac{\partial}{\partial x_\mu} + [c_\rho (x_\mu \dot{A}_\rho^c - x_\rho \dot{A}_\mu^c) - c_\mu x_\rho \dot{A}_\rho^c] \frac{\partial}{\partial A_\mu^c},$$

which is in accordance with [14].

Self-dual Yang-Mills equations and related field equations

Next, we find generators of the symmetry group admissible by self-dual Yang-Mills equations

$$F_{\mu\nu} = \frac{1}{2} \varepsilon_{\mu\nu\alpha\beta} F_{\alpha\beta}, \quad (6)$$

where $F_{\mu\nu} = A_{\mu,\nu} - A_{\nu,\mu} + [A_\mu, A_\nu]$ and $A_\mu = -\frac{i}{2} \sigma_a \dot{A}_\mu^a$. σ_a are the Pauli matrices and A_μ^a are the self-dual gauge potentials.

In fact, the solutions of equations (6) also satisfy Yang-Mills equations. They have the so-called instanton solutions [15].

Using the method described above, it can be shown that the equations (6) are invariant under the transformation group represented by the generators

$$X_\mu = \frac{\partial}{\partial x^\mu}, \quad \bar{X}_{\mu\nu} = \bar{Z}_{\mu\nu} + Y_{\mu\nu}, \quad \bar{X} = x^\mu \frac{\partial}{\partial x^\mu} - \bar{A}^\alpha \frac{\partial}{\partial A_\alpha^c}, \quad (7)$$

which form the Weyl algebra. Introducing the new variables

$$y = (x_1 + ix_2)/2, \quad \bar{y} = (x_1 - ix_2)/2, \\ z = (x_3 - ix_4)/2, \quad \bar{z} = (x_3 + ix_4)/2$$

one obtains for the equations (6) the form

$$F_{yz} = 0, \quad F_{\bar{y}\bar{z}} = 0, \quad F_{y\bar{y}} + F_{z\bar{z}} = 0. \quad (8)$$

It is known [15] that the solution of the first two equations is

$$A_y = D^{-1} D_{,y}, \quad A_{\bar{y}} = -D_{,y}^+ D^{+-1}, \\ A_z = D^{-1} D_{,z}, \quad A_{\bar{z}} = -D_{,z}^+ D^{+-1},$$

where $D \in SL(2, C)$.

The third equation for the variable $\chi = DD^+$ has the following form

$$(\chi^{-1} \cdot \chi_{,y})_{,\bar{y}} + (\chi^{-1} \cdot \chi_{,z})_{,\bar{z}} = 0, \quad (9)$$

where

$$\chi \in SL(2, C).$$

First, we search for variable χ in the form

$$\chi = q^0 I + q^\alpha \sigma_\alpha = q^\alpha \sigma_\alpha \quad (\sigma^0 \equiv I),$$

then one can get for q^α the equation

$$q_{,y\bar{y}}^\alpha + q_{,z\bar{z}}^\alpha + q^\alpha (q_{,y}^\beta q_{,\bar{y}}^\beta + q_{,z}^\beta q_{,\bar{z}}^\beta) + \\ + i\varepsilon^{\alpha\beta\gamma\delta} q^\beta (q_{,y}^\gamma q_{,\bar{y}}^\delta + q_{,z}^\gamma q_{,\bar{z}}^\delta) = 0. \quad (10)$$

In terms of the former variables x_μ , this equation has the form

$$q_{,\mu\mu}^\alpha + q^\alpha (q_{,\mu}^\beta q_{,\mu}^\beta) + 2\varepsilon^{\alpha\beta\gamma\delta} q^\beta (q_{,1}^\gamma q_{,2}^\delta - q_{,3}^\gamma q_{,4}^\delta) = 0. \quad (11)$$

The symmetry group generators of this equation form a space L_{15} :

$$X_\mu = \frac{\partial}{\partial x_\mu}, \\ X_5 = x^\mu \frac{\partial}{\partial x^\mu}, \\ X_6 = x_1 \frac{\partial}{\partial x_2} - x_2 \frac{\partial}{\partial x_1}, \\ X_7 = x_3 \frac{\partial}{\partial x_4} - x_4 \frac{\partial}{\partial x_3}, \\ X_8 = x_1 \frac{\partial}{\partial x_3} - x_3 \frac{\partial}{\partial x_1} - x_2 \frac{\partial}{\partial x_4} + x_4 \frac{\partial}{\partial x_2}, \quad (12)$$

$$X_9 = x_1 \frac{\partial}{\partial x_4} - x_4 \frac{\partial}{\partial x_1} + x_2 \frac{\partial}{\partial x_3} - x_3 \frac{\partial}{\partial x_2},$$

$$X_{\alpha\beta} = q_\alpha \frac{\partial}{\partial q_\beta} - q_\beta \frac{\partial}{\partial q_\alpha}.$$

Compared to generators (7), the Lorentz-invariance gets lost here. Generators $Z_{\mu\nu}$ in (7) correspond to generators $X_6 - X_9$ in (12). A new symmetry appears — dilatation X_5 .

Secondly, in equation (9) variable χ can be parametrised by the use of Poincaré coordinates [15]:

$$\chi = \frac{1}{u^1} \begin{pmatrix} 1 & u^3 \\ u^2 & (u^1)^2 + u^2 u^3 \end{pmatrix},$$

where $u^2 \equiv \overline{u^3}$. Then one obtains the system of Yang equations [16]:

$$\begin{aligned} u^1(u^1_{,yy} + u^1_{,zz}) - u^1_{,y}u^1_{,y} - u^1_{,z}u^1_{,z} + u^2_{,y}u^3_{,y} + u^2_{,z}u^3_{,z} &= 0, \\ u^1(u^2_{,yy} + u^2_{,zz}) - 2u^2_{,y}u^1_{,y} - 2u^2_{,z}u^1_{,z} &= 0, \\ u^1(u^3_{,yy} + u^3_{,zz}) - 2u^1_{,y}u^3_{,y} - 2u^1_{,z}u^3_{,z} &= 0, \end{aligned} \quad (13)$$

which is equivalent to the system of Ernst equations [17-19]. The symmetry group of the system (13) is the following:

$$\begin{aligned} X_\mu &= \frac{\partial}{\partial x_\mu}, \\ X_5 &= x^\mu \frac{\partial}{\partial x^\mu}, \\ X_6 &= x_1 \frac{\partial}{\partial x_2} - x_2 \frac{\partial}{\partial x_1}, \\ X_7 &= x_3 \frac{\partial}{\partial x_4} - x_4 \frac{\partial}{\partial x_3}, \\ X_8 &= x_1 \frac{\partial}{\partial x_3} - x_3 \frac{\partial}{\partial x_1} - x_2 \frac{\partial}{\partial x_4} + x_4 \frac{\partial}{\partial x_2}, \\ X_9 &= x_1 \frac{\partial}{\partial x_4} - x_4 \frac{\partial}{\partial x_1} + x_2 \frac{\partial}{\partial x_3} - x_3 \frac{\partial}{\partial x_2}, \\ X_{10} &= u^1 \frac{\partial}{\partial u^1} + u^2 \frac{\partial}{\partial u^2}, \\ X_{11} &= u^1 \frac{\partial}{\partial u^1} + u^3 \frac{\partial}{\partial u^3}, \\ X_{12} &= a(\bar{y}, \bar{z}) \frac{\partial}{\partial u_2}, \end{aligned} \quad (14)$$

$$X_{13}=b(y,z)\frac{\partial}{\partial u_3},$$

$$X_{14}=f(\bar{y},\bar{z})\left[u^1u^2\frac{\partial}{\partial u^1}+(u^2)^2\frac{\partial}{\partial u^2}-(u^1)^2\frac{\partial}{\partial u^3}\right],$$

$$X_{15}=g(y,z)\left[u^1u^3\frac{\partial}{\partial u^1}-(u^1)^2\frac{\partial}{\partial u^2}+(u^3)^2\frac{\partial}{\partial u^3}\right],$$

where a, b and f, g are arbitrary functions. Therefore the corresponding space is L_∞ .

Comparison with the symmetry generators (12) shows that subalgebra $L_9 \equiv \{X_1 - X_9\}$ of spatial transformations is the same. The other generators connected with fields are different. Initial field variables can be expressed as

$$A_y^a = \frac{1}{u^1} [iu_{,y}^2, u_{,y}^2, -iu_{,y}^4], \quad A_z^a = \frac{1}{u^1} [iu_{,z}^3, u_{,z}^3, -iu_{,z}^4].$$

Supposing that fields $q^\alpha(x)$ in (10) depend only on two coordinates

$$q^\alpha = q^\alpha(x_1, x_3),$$

the equation will be

$$q_{,11}^\alpha + q_{,33}^\alpha + q^\alpha (q_{,1}^\beta q_{,1}^\beta + q_{,3}^\beta q_{,3}^\beta) = 0, \quad (15)$$

which in terms of new variables $w = \frac{1}{2}(x_1 + ix_3)$, \bar{w} means that

$$q_{,ww}^\alpha + q^\alpha (q_{,w}^\beta q_{,w}^\beta) = 0. \quad (16)$$

Imposing the supplementary condition $q^2 = 1$, one can obtain the non-linear $O(1,3)$ σ -model [15]. The symmetry group of equations (15) is the following —

$$\begin{aligned} X_1 &= f(x_1 + ix_3) \left(\frac{\partial}{\partial x_1} - i \frac{\partial}{\partial x_3} \right), \\ X_2 &= g(x_1 - ix_3) \left(\frac{\partial}{\partial x_1} + i \frac{\partial}{\partial x_3} \right), \end{aligned} \quad (17)$$

$$X_{\alpha\beta} = q_\alpha \frac{\partial}{\partial q_\beta} - q_\beta \frac{\partial}{\partial q_\alpha},$$

where f and g are arbitrary functions.

Generators X_1 and X_2 include shifts, rotations and dilatations on the $(1,3)$ -plane. Indeed, taking $f=g=1/2$ and $f=-g=i/2$, one gets the shift generators $X_1 + X_2 = \frac{\partial}{\partial x_1}$ and $X_1 - X_2 = \frac{\partial}{\partial x_3}$, respectively.

Taking $f(w) = iw$, $g(\bar{w}) = -i\bar{w}$ and $f(w) = w$, $g(\bar{w}) = \bar{w}$, one obtains the rotation $X_1 + X_2 = x_1 \frac{\partial}{\partial x_3} - x_3 \frac{\partial}{\partial x_1}$ and the dilatation $X_1 + X_2 = x_1 \frac{\partial}{\partial x_1} + x_3 \frac{\partial}{\partial x_3}$, respectively.

In variables w and \bar{w} X_1 and X_2 take the forms $X_1=f(w) \frac{\partial}{\partial w}$ and $X_2=g(\bar{w}) \frac{\partial}{\partial \bar{w}}$.

Conserved quantities

Here we proceed with the investigation of the considered equations, trying to find some conserved quantities obtainable by means of the Noether theorem [20]. The conserved vector K_r^μ satisfies the condition

$$D_\mu K_r^\mu \equiv \mathcal{L} D_\mu \xi_r^\mu + \xi_r^\mu \frac{\partial \mathcal{L}}{\partial x^\mu} + \eta_r^\alpha \frac{\partial \mathcal{L}}{\partial u^\alpha} + \xi_{\mu r}^\alpha \frac{\partial \mathcal{L}}{\partial u^\alpha_{,\mu}} + \dots = 0. \quad (18)$$

Such a vector can be expressed in the form

$$K_r^\mu = (\eta_r^\alpha - u^\alpha_{,\nu} \xi_r^\nu) \frac{\partial \mathcal{L}}{\partial u^\alpha} + \mathcal{L} \xi_r^\mu \quad (r=1, \dots, N). \quad (19)$$

Here only the equations up to the second order are considered. ξ_r^ν and η_r^α are the coefficients in the r -th generator X_r .

$$D_\mu \equiv \frac{\partial}{\partial x^\mu} + u^\alpha_{,\mu} \frac{\partial}{\partial u^\alpha}$$

is the covariant derivative operator.

According to the Noether theorem, the action functional is required to be invariant with respect to an N -parameter continuous group of transformations G with a given system of generators $\{X_r\}$. We want to satisfy a weaker condition, i.e. the equation of motion being invariant with respect to a certain group of transformations. In principle, a generator, admitted by the equation of motion, need not leave its action invariant. For this reason one must either check the validity of equation (18) for every generator X_r , or use a generalization of the Noether theorem [8] which considers the invariance properties of extremal values of the variational integrals only, instead of regarding all their admissible values.

In case of Yang-Mills equations the knowledge of shift generators $X_\mu = \delta_{\mu\nu} \frac{\partial}{\partial x_\nu}$ of symmetry group (5) provides us with a possibility of finding obvious conserved quantities — energy and momentum. Indeed, one can write immediately

$$D_\mu K_\nu^\mu = 0$$

and

$$K_{\mu\nu} = \overset{c}{F}_{\rho\mu} \overset{c}{F}_{\rho\nu} - \frac{1}{4} \delta_{\mu\nu} \overset{c}{F}_{\rho\sigma} \overset{c}{F}_{\rho\sigma}. \quad (20)$$

This is the energy-momentum tensor, where the energy density is

$$H = -K_{11} \quad (x^1 \equiv t).$$

There is also a conserved quantity-angular momentum tensor

$$K_{\alpha\beta}^\mu = x_\beta K_\alpha^\mu - x_\alpha K_\beta^\mu + \overset{c}{F}_{\mu\alpha} \overset{c}{A}_\beta - \overset{c}{F}_{\mu\beta} \overset{c}{A}_\alpha \quad (21)$$

for the rotation generator

$$X_{\alpha\beta} = \xi_{\alpha\beta}^{\mu} \frac{\partial}{\partial x^{\mu}} + \eta_{\alpha\beta}^{\mu} \frac{\partial}{\partial A_{\mu}^c},$$

where

$$\xi_{\alpha\beta}^{\mu} = \delta_{\mu\alpha} x_{\beta} - \delta_{\mu\beta} x_{\alpha},$$

$$\eta_{\alpha\beta}^{\mu} = \delta_{\mu\alpha} A_{\beta}^c - \delta_{\mu\beta} A_{\alpha}^c.$$

For the Yang equations (13) the Lagrangian is [15] —

$$\begin{aligned} \mathfrak{L} &= \frac{1}{2(u^1)^2} (u^1_{,y} u^1_{,y} + u^1_{,z} u^1_{,z} + u^2_{,y} u^3_{,y} + u^2_{,z} u^3_{,z}) = \\ &= \frac{1}{2(u^1)^2} \{ u^1_{,\mu} u^1_{,\mu} + u^2_{,\mu} u^3_{,\mu} + i([u^2, u^3]_{,12} + [u^2, u^3]_{,43}) \}, \end{aligned} \quad (22)$$

where $[u^2, u^3]_{,ik} = u^2_{,i} u^3_{,k} - u^2_{,k} u^3_{,i}$ are the Poisson brackets. Condition $D_{\mu} K^{\mu}_r = 0$ holds only for conserved quantities which correspond to the generators $X_v, X_6, X_7, X_8, X_9, X_{12}, X_{13}$. For example,

$$\begin{aligned} K^{\mu}_v &= \frac{1}{2(u^1)^2} \{ (u^1_{,\rho} u^1_{,\rho} + u^2_{,\rho} u^3_{,\rho}) \delta_{\mu v} - 2u^1_{,\mu} u^1_{,v} - u^2_{,\mu} u^3_{,v} - \\ &- u^2_{,v} u^3_{,\mu} + i\delta_{\mu v} ([u^2, u^3]_{,12} + [u^2, u^3]_{,43}) + i\delta_{\mu 1} [u^2, u^3]_{,2v} + \\ &+ i\delta_{\mu 2} [u^2, u^3]_{,v1} + i\delta_{\mu 3} [u^2, u^3]_{,v4} + i\delta_{\mu 4} [u^2, u^3]_{,3v} \}. \end{aligned} \quad (23)$$

This is the energy-momentum tensor.

$$\begin{aligned} K^1_6 &= \frac{1}{2(u^1)^2} \left\{ x_2 \left(u^1_{,1} u^1_{,1} + u^2_{,1} u^3_{,1} - \sum_{i=2,3,4} (u^1_{,i} u^1_{,i} + u^2_{,i} u^3_{,i}) + \right. \right. \\ &\left. \left. + i[u^2, u^3]_{,34} \right) - x_1 (2u^1_{,1} u^1_{,2} + \{u^2, u^3\}_{,12}) \right\}, \end{aligned} \quad (24)$$

where $\{u^2, u^3\}_{,ik} = u^2_{,i} u^3_{,k} + u^2_{,k} u^3_{,i}$;

$$K^{\mu}_{13} = \frac{b(y, z)}{2(u^1)^2} [u^2_{,\mu} - i(u^2_{,2} \delta_{\mu 1} - u^2_{,4} \delta_{\mu 2} - u^2_{,4} \delta_{\mu 3} + u^2_{,3} \delta_{\mu 4})]. \quad (25)$$

The other conserved quantities $K^2_6, K^3_6, K^4_6, K^{\mu}_7, K^{\mu}_8, K^{\mu}_9, K^{\mu}_{12}$ can be obtained analogously by means of equation (19).

Also, for the nonlinear σ -model equation (15) with the Lagrangian [15]

$$\mathfrak{L} = \frac{1}{2} \partial_{\mu} q^{\alpha} \partial_{\mu} q^{\alpha} \quad (\mu = 1, 2) \quad (26)$$

and with the symmetry group (17) one can obtain conserved vectors $K^{\mu}_1, K^{\mu}_2, K^{\mu}_3$, where

$$X_3 = \frac{1}{2} \varepsilon_{\alpha\beta} X_{\alpha\beta} \quad (\varepsilon_{\alpha\beta} = -\varepsilon_{\beta\alpha}). \quad (27)$$

Indeed, $D_{\mu} K^{\mu}_3 = 0$ and

$$K_3^\mu = \varepsilon_{\alpha\nu} q_\nu q_{\alpha,\mu} \quad (28)$$

is the conserved spin-vector.

Conclusions and acknowledgements

The generators of symmetry groups and several conserved quantities have been found for Yang-Mills equations (at the Lorentz gauge), self-dual Yang-Mills equations and the equations deduced from the latter (Yang equations, nonlinear σ -model equations, etc.).

Not every generator of an equation is associated with a conserved quantity. For obtaining the conserved currents, the Lagrangian is needed, but, unfortunately, not every equation is provided by the Lagrangian.

In such cases the question of existence and finding the systems of conserved quantities remains open.

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Academy of Sciences of the Estonian SSR,
Institute of Physics

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V. ROSENHAUS, K. KIIRANEN

**YANG-MILLSI VÖRRANDITE JA OMADUAALSE YANG-MILLSI VÖRRANDI
SÜMMEETRIARÜHMADEST**

Artikkel käsitleb Yang-Millsi võrrandite ja omaduaalse Yang-Millsi võrrandi mõnede kujude poolt lubatud punkteisenduste rühmade leidmise probleemi neljamõõtmelises eukleidilises ruumis. On leitud nende rühmade baasgeneraatorid, samuti mõned jäävad suurused.

В. РОЗЕНГАУЗ, К. КИИРАНЕН

**О ГРУППАХ СИММЕТРИИ УРАВНЕНИЙ ЯНГА—МИЛЛСА
И САМОДУАЛЬНОГО УРАВНЕНИЯ ЯНГА—МИЛЛСА**

Рассматривается задача отыскания групп точечных преобразований, допускаемых уравнением Янга—Миллса и некоторыми видами уравнений Янга—Миллса самодуальных полей в четырехмерном евклидовом пространстве. Найден явный вид базисных генераторов данных групп. Получены с помощью теоремы Нётер некоторые сохраняющиеся величины.